

# Perturbed Power Law Parameters

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# Perturbed Power-Law Parameters from WMAP7

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Power spectrum

Observations

Cosmological Parameters

Inflation

Gaussian, Adiabatic, nearly Scale Invariant perturbations

Power Law Inflation Models - uniform accelerated approximation

# Power Law Inflation

Power law model of inflation  $a(t) = t^p$   $p > 1$

$$\epsilon = -\delta = \frac{1}{p}$$

i. uniformly accelerated expansion  $-\frac{\dot{H}}{H^2} = \text{const}$

ii.  $\delta\phi$ , massless scalar field perturbations,  $m_{\text{eff}}^2 = 0$

Scalar and tensor perturbations spectra of identical shape with const  $n_s$  and  $n_t$  ( $n_s = n_t + 1$ )

Standard slow-roll inflation

Scale invariant spectrum,  $\mathcal{P}(k) \propto k^{(n-1)}$

$n_s = 1 \rightarrow$  Scale free Harrison - Zeldovich spectrum

$$\text{Spectral Index, } n = 1 + \frac{d \ln P(k)}{d \ln k}$$

Relative distribution of power on various scales

$$\text{Running of the Spectral Index } \frac{dn}{d \ln k}$$

Deviation of the primordial power spectrum from power law

Models with soft departures from scale free spectra

Hubble parameter  $H(\phi)$

$$V(\phi) = \frac{3m_{\mathcal{P}}^2}{8\pi} H^2(\phi) \left[ 1 - \frac{m_{\mathcal{P}}^2}{12\pi} \left( \frac{\partial \ln H}{\partial \phi} \right)^2 \right] \quad (1)$$

The field equation for the modes of the inflaton perturbation,  $\delta\phi_k$

$$\delta\phi_k'' + 2\frac{a'}{a}\delta\phi_k' + (k^2 + a^2 m_{\text{eff}}^2)\delta\phi_k = 0, \quad (2)$$

$$a^2 m_{\text{eff}}^2 = \frac{a''}{a} - \frac{z''}{z} \text{ and } z = a^2 \frac{\phi'}{a'}.$$

Effective mass

$$\frac{m_{eff}^2}{H^2} = -(\epsilon + \delta)(\delta + 3) + \frac{\dot{\epsilon}}{H} - \frac{\dot{\delta}}{H} \quad (3)$$

Slow-roll parameters

$$\epsilon = \frac{m_{\mathcal{P}}^2}{4\pi} \left( \frac{H_{\phi}}{H} \right)^2 \quad \text{and} \quad \delta = -\frac{m_{\mathcal{P}}^2}{4\pi} \frac{H_{\phi\phi}}{H} \quad (4)$$

$\frac{H_{\phi}}{H} = \left( \frac{\partial \ln H}{\partial \phi} \right)$  is a constant for the power law models with uniform acceleration  $\Rightarrow \frac{m_{eff}^2}{H^2}$  will be zero

# Power Law Model

Scalar and tensor perturbation spectra - parametrized by a common spectral index  $\nu = n_s - 1 = n_t$

$$\nu = \frac{-2\epsilon}{1 - \epsilon} = \frac{-4 \left( \frac{\partial \ln H}{\partial \phi} \right)^2}{\left[ 1 - 2 \left( \frac{\partial \ln H}{\partial \phi} \right)^2 \right]} \quad (5)$$

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$$a(\eta) = \left( \frac{\eta}{\eta_*} \right)^{(1/2-\mu)} \quad \mu = \frac{3}{2} + \frac{\epsilon(\eta)}{1-\epsilon(\eta)}$$

$$H^2(\phi) = \frac{(\mu - 1/2)^2}{\eta^2 a^2(\eta)}, \quad \frac{z''}{z} = \frac{1}{\eta^2} \left[ \mu_S^2 - \frac{1}{4} \right] \quad (6)$$

$$\mu_S^2 = \mu^2 - (\mu - 1/2)^2 \frac{m_{\text{eff}}^2}{H^2}$$

# Perturbed Power Law

Small deviations from uniform acceleration -

→  $\partial^2 \ln H / \partial \phi^2$  perturbation

Scalar and Tensor spectra are perturbed from the scale free form  
At the leading order maintain a constant difference between their spectral indices

$$\frac{m_{\text{eff}}^2}{H^2} \simeq \text{const} \text{ (weakly depend on } \eta \text{)}$$

PPL models -  $n_s \neq n_t + 1$

$$n_s - 1 = \nu_s$$

$$n_t = \nu_t$$

$$\nu_{st} = \nu_s - \nu_t$$

$$\nu_s(k) = \frac{-2\epsilon(\eta_*)}{1 - \epsilon(\eta_*)} + \frac{2\chi(\eta_*)}{1 - \epsilon(\eta_*)} + \mathcal{O}(\chi^2) \simeq \frac{-4 \left[ \left( \frac{\partial \ln H}{\partial \phi} \right)^2 + \left( \frac{\partial^2 \ln H}{\partial \phi^2} \right) \right]}{\left[ 1 - 2 \left( \frac{\partial \ln H}{\partial \phi} \right)^2 \right]} \quad (7)$$

$$\chi = \epsilon + \delta = -2 \left( \frac{\partial^2 \ln H}{\partial \phi^2} \right) \quad (8)$$

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$$\chi = \epsilon + \delta = -2 \left( \frac{\partial^2 \ln H}{\partial \phi^2} \right) \quad (8)$$

$$\nu_{st}(k) = \frac{2\chi(\eta_*)}{1 - \epsilon(\eta_*)} = \frac{-4 \left( \frac{\partial^2 \ln H}{\partial \phi^2} \right)}{\left[ 1 - 2 \left( \frac{\partial \ln H}{\partial \phi} \right)^2 \right]} \quad (9)$$

Scalar perturbation spectrum

$$\mathcal{P}_s(k) = \frac{\mathcal{A}(\mu_T, \mu_S)}{4\pi} \left( \frac{H_*}{2\pi} \right)^2 \frac{1}{\epsilon(\eta_*)} \quad (10)$$

Tensor perturbation spectrum

$$\mathcal{P}_t(k) = \frac{8 \mathcal{A}(\mu_T, \mu_T)}{2\pi} \left( \frac{H_*}{2\pi} \right)^2 \quad (11)$$

$$\mathcal{A}(x, y) = \frac{4^y \Gamma^2(y)}{(x-1/2)^{2y-1}}$$

Standard slow-roll approximation

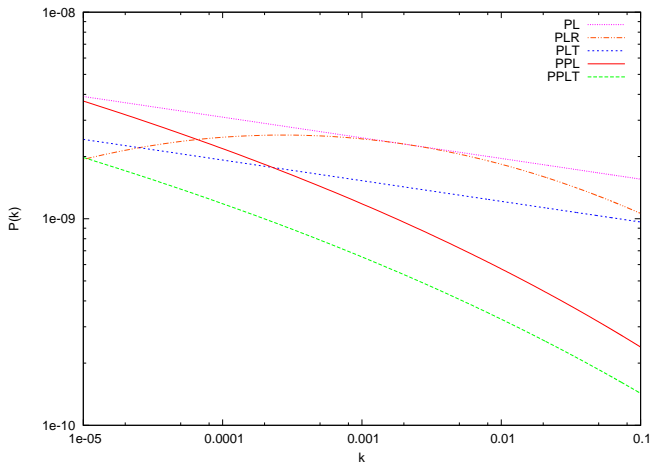
$$\mu_T = \mu_S = \mu$$

PPL approach can be used to study models with comparatively larger  $\epsilon$  and  $\delta$  with the condition that  $\chi = \epsilon + \delta$  is small ( $\chi$  is zero for PL case)

$$\mathcal{P}_s(k) \quad \text{and} \quad \mathcal{P}_t(k)$$

**Running** of their spectral indices also

# PPL Spectra



PL spectrum for a model with  $n_s = 0.9$ , PLR spectrum with  $n_s = 0.9$  and  $n_{run} = -0.05$  and PLT spectrum with  $n_t = -0.1$ . The scalar and tensor PPL spectra are plotted with  $\nu_s = -0.1$  ( $n_s = 0.9$ ) and  $\nu_{st} = -0.15$



- $\mathcal{P}_s(k)$  and  $\mathcal{P}_t(k)$

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T. Souradeep, J. R. Bond and L. Knox *et al*, 1998

Prospects for measuring Inflation parameters with the CMB : Proceedings of COSMO-97 - International workshop on Particle Physics and Early Universe, [arXiv:astro-ph/9802262].

# Parameter Estimation with PPL

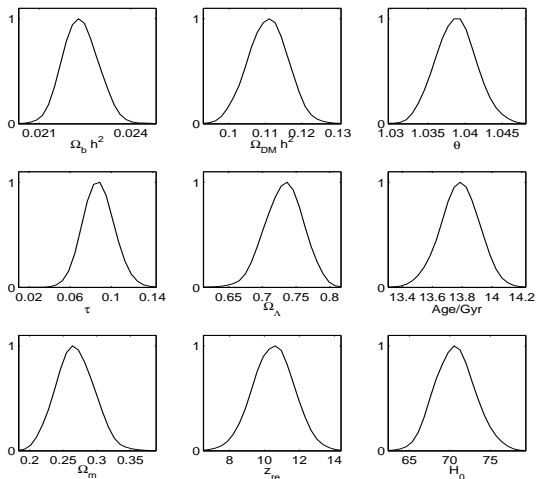
Markov Chain Monte Carlo sampling of the parameter space

Constraints on the inflationary parameters and the various background cosmological parameters

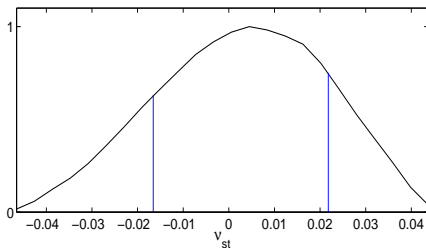
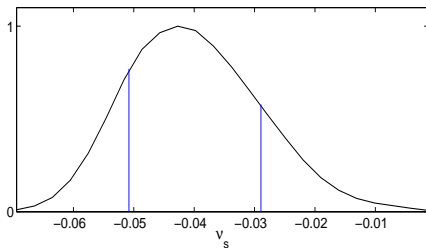
Parameter	Lower Limit	Upper Limit
$\Omega_b h^2$	0.005	0.1
$\Omega_c h^2$	0.01	0.99
$\theta$	0.5	10
$\tau$	0.01	0.8
$\nu_s$	-0.15	0.0
$\nu_{st}$	-0.06	0.06
$\log[10^{10} A_s]$	2.0	4.2

Priors for the background cosmological parameters and the PPL inflationary parameters

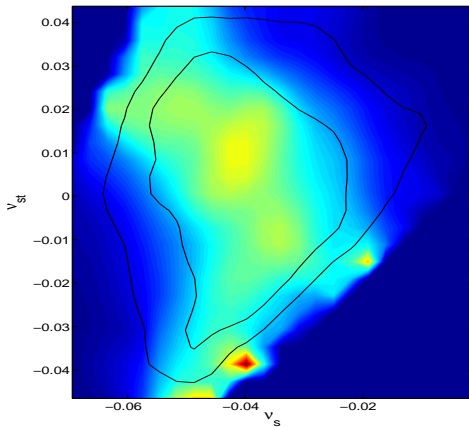
# Background Cosmological Parameters



1D marginalized posterior distribution of the background cosmological parameters obtained by PPL fitting to WMAP7 data

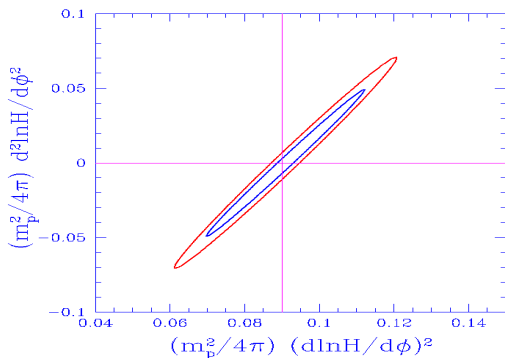


1D marginalized posterior distribution of the PPL inflationary parameters  $\nu_s$  and  $\nu_{st}$  obtained by fitting to WMAP7 data. The 95% marginalised limit on the parameters are shown by the vertical lines.

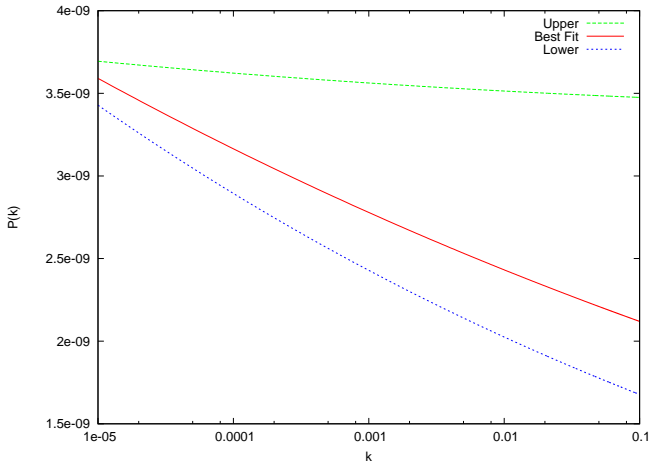


The joint 2D marginalized posterior distribution of  $\nu_s$  and  $\nu_{st}$

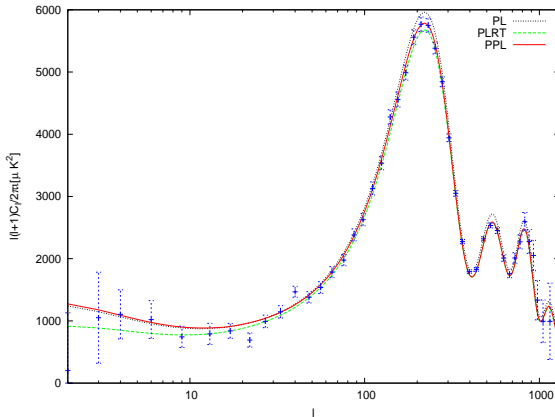
## Forecast



The expected 1- $\sigma$  error ellipses in the  $(\partial \ln H/\partial\phi)^2 - (\partial^2 \ln H/\partial\phi^2)$



Ranges of possible deviation from PL spectrum with in 95% CL. Best fit PPL  $\mathcal{P}(k)$  (red, solid) and the PPL spectra with upper ( $\nu_s = -0.03$  and  $\nu_{st} = 0.02$ ) and with lower ( $\nu_s = -0.05$  and  $\nu_{st} = -0.02$ ) limit of 95% CL of  $\nu_s$  and  $\nu_{st}$



Best fit  $C_\ell$ s for power-law (PL) (black, dotted), power law with tensor and running of scalar spectral index (PLRT) (green, dashed), power law with perturbative method (PPL) (red, solid)

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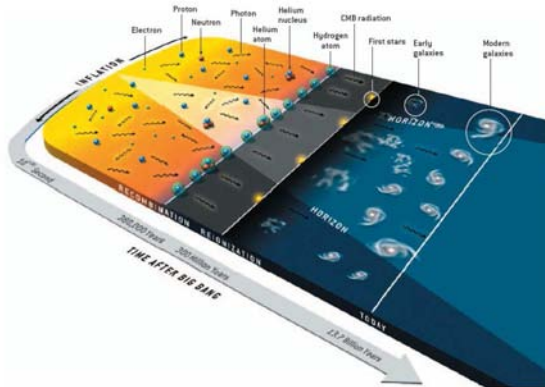
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- With PPL just the two parameters  $\nu_s$  and  $\nu_{st}$  will capture all these.
- PPL spectra takes care of the running of tensor spectral index also
- In contrast to the slow roll expansion method, for which  $\epsilon$  and  $\delta$  have to be small, PPL can look also at models with comparatively larger  $\epsilon$  and  $\delta$  with the condition that  $(\epsilon + \delta)$  is small.

## Perturbed Power Law

How the general models of inflation are best studied as departures from the power law model.

The probable values of the **deviation parameter**  $\nu_{st}$   
Contained between  $-0.04 < \nu_{st} < 0.04$ .

# Thank You !





Model	No. of extra Parameters	$\chi_{eff}^2$
<b>PPL</b>	1	7474.96
PL	0	7475.14
PLT	2	7475.00
PLTC	1	7475.07
PLR	1	7473.74
PLRT	3	7473.58
PLRTC	2	7473.61

The best fit likelihood  $\chi_{eff}^2$  values for the perturbed power law (PPL), power law (PL), power law with tensor (PLT), power law with tensor obeying the consistency relation  $n_t = -r/8$  (PLTC), power law with running of scalar spectral index (PLR), power law with running and tensor (PLRT), power law with running and tensor obeying the consistency relation (PLRTC) models.