

Indo-UK Seminar

Cosmological viability of $f(R)$ theory of gravity

by

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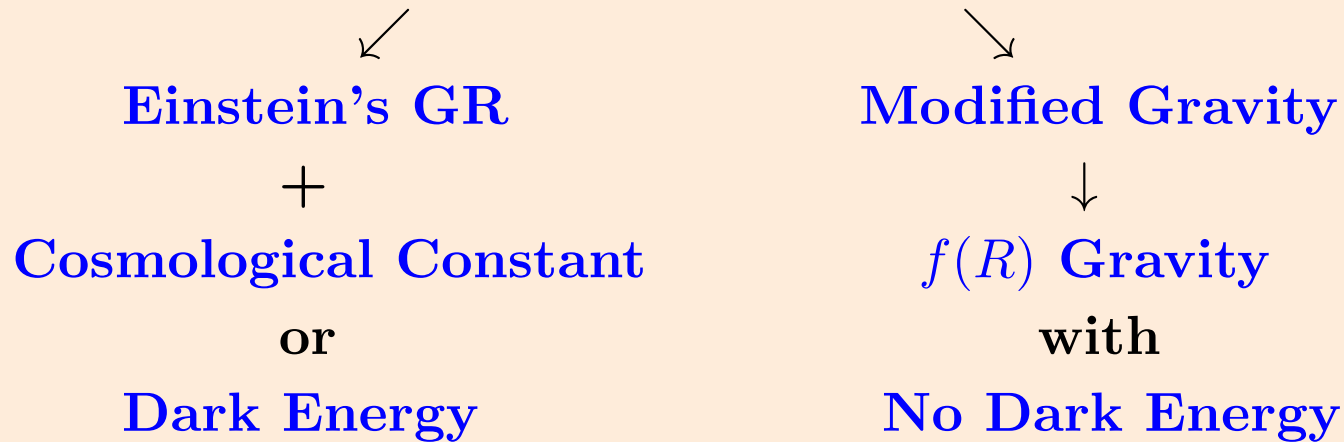
IUCAA, Pune

Based on [arXiv:1106.6353](https://arxiv.org/abs/1106.6353)

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Late time accelerated expansion



- From the cosmological dynamics are $f(R)$ theories distinguishable from dark energy ?
- Can one design an $f(R)$ theory for each DE model ?
- In particular is it possible to design an $f(R)$ theory which behaves as dark energy with constant w .

Alternative to dark energy $\rightarrow f(R)$ theory

- May be Einstein's GR is not correct on large scales !

$$\mathcal{S} = \frac{1}{16\pi G} \int (R + f(R)) \sqrt{-g} d^4x + \mathcal{S}_m$$

where \mathcal{S}_m is the action for the matter content

S. Capozziello *et. al.* (2003) & S. Carroll *et. al.* (2004)
suggested that

$$\mathcal{S} = \frac{1}{16\pi G} \int d^4x \sqrt{-g} \left(R - \frac{\mu^4}{R} \right) + \mathcal{S}_m$$

can explain the accelerated expansion without dark energy!

A deeper look at the $f(R)$ -theory

$$\mathcal{S} = \frac{1}{16\pi G} \int (R + f(R)) \sqrt{-g} d^4x + \mathcal{S}_m$$

Modified Einstein's Equation

$$(1 + f_R)G^\mu{}_\nu - g^{\mu\alpha} f_{R,\alpha;\nu} + \left(\frac{2\Box f_R - (f - Rf_R)}{2} \right) \delta^\mu{}_\nu = (8\pi G) T_{(m)\nu}{}^\mu$$

where $f_R \equiv \frac{df}{dR}$

Trace of the above equation

$$\Box f_R + \left(\frac{Rf_R - (2f + R)}{3} \right) = \left(\frac{8\pi G}{3} \right) T$$

The 'hidden' scalar field in $f(R)$ theory

$$\phi \equiv 1 + f_R$$

Field equations in $f(R)$ theory becomes

$$G_{\mu\nu} = \left(\frac{8\pi G}{\phi} \right) T_{\mu\nu} + \frac{1}{\phi} (\phi_{;\mu;\nu} - g_{\mu\nu} \square \phi) + \left(\frac{V(\phi)}{2} \right) g_{\mu\nu}$$

$$\square \phi + \left(\frac{1}{3} \right) \left(\phi \left(\frac{dV}{d\phi} \right) - 2V(\phi) \right) = \left(\frac{8\pi G}{3} \right) T$$

where $V(\phi) \equiv f - Rf_R$

These equations are *identical* to the one in Brans-Dicke theory

$$\mathcal{S} = \frac{1}{16\pi G} \int \left(\phi R - \omega \left(\frac{\partial_\mu \phi \partial^\mu \phi}{\phi} \right) + V(\phi) \right) \sqrt{-g} d^4x + \mathcal{S}_m$$

but with the BD parameter $\omega = 0$.

$f(R)$ theory \iff Brans-Dicke theory with $\omega = 0$

- $f(R)$ theory of gravity are equivalent to Brans-Dicke theory with the parameter $\omega = 0$

$$\mathcal{S} = \frac{1}{16\pi G} \int (R + f(R)) \sqrt{-g} d^4x + \mathcal{S}_m$$

\iff

$$\mathcal{S} = \frac{1}{16\pi G} \int (\phi R + V(\phi)) \sqrt{-g} d^4x + \mathcal{S}_m$$

- This equivalence is strictly valid when $f_{RR} \neq 0$

The designer $f(R)$ theory

- We want to design an $f(R)$ theory for which $a(t)$ satisfies the following equation

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \left(\frac{8\pi G}{3}\right) \left[\rho_{m0} a^{-3} + \rho_{d0} a^{-3(1+w)}\right]$$

The reconstruction equations:

$$f'_R + 3 \left(\frac{1 + 3w\Omega_d(a)}{2a}\right) f_R - \frac{f}{6aH^2} = -\frac{\Omega_d(a)}{a}$$

$$f''_R - 3 \left(\frac{1 + w\Omega_d(a)}{2a}\right) f'_R - 3 \left(\frac{1 + w\Omega_d(a)}{a^2}\right) f_R = \frac{3(1+w)\Omega_d(a)}{a^2}$$

where $\Omega_d(a) = \frac{\rho_{d0} a^{-3(1+w)}}{\rho_{m0} a^{-3} + \rho_{d0} a^{-3(1+w)}}$

Approximate $f(R)$ at $a \ll 1$

In the limit $a \ll 1$, the reconstruction equations reduces to

$$f'_R + \frac{3}{2} \left(\frac{f_R}{a} \right) - \frac{f}{6 a H^2} = - \left(\frac{\Omega_{d0}}{\Omega_{m0}} \right) a^{-(1+3w)}$$

$$f''_R - \frac{3}{2} \left(\frac{f'_R}{a} \right) - 3 \left(\frac{f_R}{a^2} \right) = \left(\frac{\Omega_{d0}}{\Omega_{m0}} \right) a^{-(2+3w)}$$

On solving gives:

$$f_R(a) = C_1 a^3$$

where C_1 is the constant of integration.

The Constant C_1

The consistency requirement implies that

$$C_1 = 3(1+w) \left(\frac{\Omega_{d0}}{\Omega_{m0}} \right) a^{-3(1+w)}$$

But C_1 is the constant of integration !!

Therefore, for an $f(R)$ theory to mimic DE with constant w

- The only allowed value of w is $w = -1$
- This makes $C_1 = 0$
- Consequently $f_R = 0$, $f'_R = 0$, $f''_R = 0$, and so and so forth

The only possible solution is

$$f(R) = \text{constant}$$

Conclusion

If $f(R)$ theory with action

$$\mathcal{S} = \frac{1}{16\pi G} \int (R + f(R)) \sqrt{-g} d^4x + \mathcal{S}_m$$

is to drive the the accelerated expansion of the universe as if it is due to dark energy with a constant w in Einstein's GR

\Rightarrow The only allowed value of w for which it can do so is

$$w = -1$$

\Rightarrow But in this case the action reduces to

$$\mathcal{S} = \frac{1}{16\pi G} \int (R + 2\Lambda) \sqrt{-g} d^4x + \mathcal{S}_m$$

Implications

- The result doesn't mean that $f(R)$ theories are not a viable alternative to dark energy.
- In fact, it demonstrates that in $f(R)$ theories, the resulting equation of state parameter of the equivalent dark energy must necessarily be epoch dependent.

Immediate corollary to the above result

- $w(a)$ in $f(R)$ theories must necessarily be very close to -1 during the matter dominated epoch

S. Unnikrishnan, S. Thakur, T. R. Seshadri, [arXiv:1106.6353](https://arxiv.org/abs/1106.6353)

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* **Thank you** *

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