

NON-GAUSSIAN FLUCTUATIONS OF THE INFLATON AND CONSTANCY OF CORRELATIONS OF ζ OUTSIDE THE HORIZON

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- IN THE SIMPLEST COSMOLOGICAL MODELS, CURVATURE PERTURBATIONS ARE GENERATED BY FLUCTUATIONS IN THE INFLATON FIELD.
- THESE CURVATURE PERTURBATIONS THEN LEAD TO FLUCTUATIONS IN THE CMBR AND IN MATTER.
- CURVATURE PERTURBATIONS CAN BE DESCRIBED BY A VARIABLE ζ , AND CORRELATIONS OF ζ LEAD TO CORRELATIONS IN THE CMBR AND MATTER DISTRIBUTION.
- GAUSSIAN ($\langle \zeta \zeta \rangle$ NON-TRIVIAL, 4PT $\sim (2 \text{ PT})^2$, ODD=0)
NON-GAUSSIAN (ODD, OR NON-TRIVIAL EVEN)

NON-GAUSSIANITIES IN SINGLE FIELD INFLATION

- NON-GAUSSIANITIES IN THE CURVATURE PERTURBATION ζ CAN ARISE FROM
 - 1) SELF INTERACTIONS OF THE INFLATON
 - 2) NON-LINEARITIES IN COSMOLOGICAL PERTURBATION THEORY
 - 3) PREHEATING, 4) NON-VACUUM QUANTUM STATE

USUALLY THE SELF INTERACTION CONTRIBUTION IS IGNORED.

SELF INTERACTION CONTRIBUTION

FOR EXAMPLE, CONSIDER CUBIC SELF INTERACTIONS,

$\mu \phi^3$, WHERE ϕ IS THE INFLATON

- THE SELF INTERACTION CONTRIBUTION TO $\langle \hat{\zeta}^3 \rangle \sim \langle (\hat{\delta\phi})^3 \rangle$ WHICH IS PROPORTIONAL TO μ , AND THUS TO THE SLOW ROLL PARAMETER $\xi \sim V'''$

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- IS PRESUMED SMALLER THAN TERMS PROPORTIONAL TO SLOW ROLL PARAMETERS $\epsilon \sim V'^2$ AND $\eta \sim V''$, AND SO IGNORED. [MALDACENA 2003]

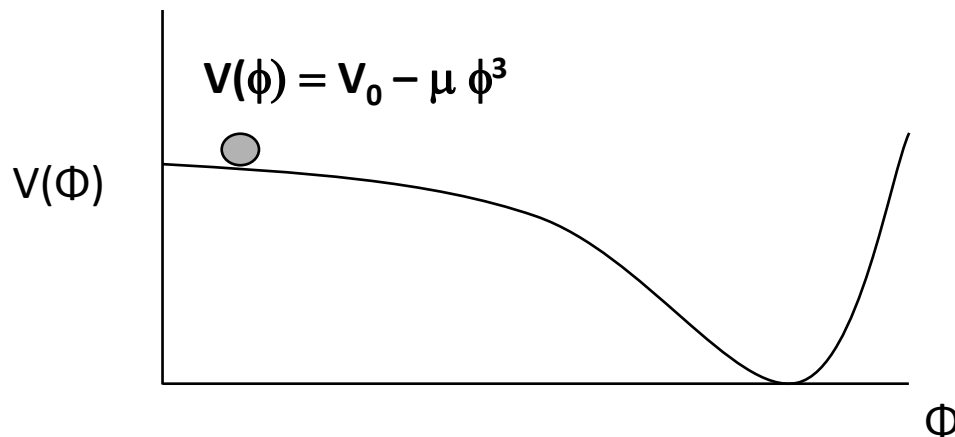
BUT

NON-GAUSSIANITIES IN SINGLE FIELD INFLATION

1. FOR MANY MODELS (NEW INFLATION, SMALL FIELD NATURAL INFLATION AND RUNNING MASS INFLATION), $\xi \gg \epsilon$.

FOR EXAMPLE, NEW INFLATION $V(\phi) = V_0 - \mu \phi^3$.

$$\epsilon \equiv \frac{1}{2} \left(\frac{V'}{V} \right)^2 M_P^2 \sim \frac{\mu^2 \phi^2}{H^4} \frac{\phi^2}{M_P^2} \quad \xi \equiv \frac{V'V'''}{V^2} M_P^4 \sim \frac{\mu^2 \phi^2}{H^4}$$



NON-GAUSSIANITIES IN SINGLE FIELD INFLATION

1. FOR MANY MODELS (NEW INFLATION, SMALL FIELD NATURAL INFLATION AND RUNNING MASS INFLATION), $\xi \gg \epsilon$.
2. THE SELF INTERACTION CONTRIBUTION IS ACTUALLY $\sim \xi N_e$. THIS IS COMPARABLE TO η .

($N_e(t)$ is the number of e-foldings since horizon exit, and increases from 0 to 60 by the end of inflation for our current horizon scale).

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THEREFORE SELF INTERACTIONS SHOULD NOT BE
IGNORED OUTRIGHT

[MALDACENA 2003 - CHAOTIC]

DO SELF INTERACTIONS IMPLY GROWTH OUTSIDE THE HORIZON ?

- BUT A CONTRIBUTION TO $\langle \hat{\zeta}(k)^3 \rangle \sim N_e(t)$ IMPLIES THAT THE 3-POINT FUNCTION IS GROWING AFTER HORIZON EXIT ! THIS IS CONTRARY TO ONE'S EXPECTATIONS.

IS THE CALCULATION OF THE INFLATON SELF INTERACTION
CONTRIBUTION TO $\langle \hat{\zeta}^3 \rangle$ INCORRECT?

- IT HAS BEEN DONE INDEPENDENTLY BY FALK, RR, SREDNICKI (1993), ZALDARRIAGA (2004), BERNARDEAU, BRUNIER, UZAN (2004) AND SEERY, MALIK, LYTH(2008) IN DIFFERENT CONTEXTS.

DO SELF INTERACTIONS IMPLY GROWTH OUTSIDE THE HORIZON ?

- BUT A CONTRIBUTION TO $\langle \hat{\zeta}(\mathbf{k})^3 \rangle \sim N_e(t)$ IMPLIES THAT THE 3-POINT FUNCTION IS GROWING AFTER HORIZON EXIT ! THIS IS CONTRARY TO ONE'S EXPECTATIONS.
- IT HAS BEEN DONE INDEPENDENTLY BY FALK ET AL (1993), ZALDARRIAGA (2004), BERNARDEAU ET AL (2004) AND SEERY ET AL (2008) IN DIFFERENT CONTEXTS.
- LET US RE-EXAMINE THE ARGUMENT THAT CORRELATIONS OF THE CURVATURE PERTURBATION $\zeta(\mathbf{k})$ ARE CONSTANT OUTSIDE THE HORIZON.

- $\zeta(k)$, THE CURVATURE PERTURBATION ON UNIFORM DENSITY SLICES, WAS FIRST INTRODUCED BY BARDEEN, STEINHARDT, TURNER
- GENERALISED BY SALOPEK, BOND TO STUDY NON-GAUSSIANITIES IN THE CURVATURE PERTURBATION
- IN THE LITERATURE, THE FUNCTION $\zeta(k)$ IS SHOWN TO BE CONSTANT OUTSIDE THE HORIZON.

[SALOPEK, BOND; LYTH, WANDS; RIGOPOULOS, SHELLARD; MALIK, WANDS; LYTH, MALIK, SASAKI; LANGLOIS, VERNIZZI]

- IN THE LITERATURE, THE FUNCTION $\zeta(k)$ IS SHOWN TO BE CONSTANT OUTSIDE THE HORIZON.

CLASSICAL

WHAT DOES THIS IMPLY FOR QUANTUM n-POINT
FUNCTIONS OF $\hat{\zeta}(k)$?

- IN THE LITERATURE, THE FUNCTION $\zeta(k)$ IS SHOWN TO BE CONSTANT OUTSIDE THE HORIZON. **CLASSICAL**

WHAT DOES THIS IMPLY FOR QUANTUM n-POINT FUNCTIONS OF $\hat{\zeta}(k)$?

THE 2-POINT FUNCTION

$$\langle \hat{\zeta}(k_1) \hat{\zeta}(k_2) \rangle = (2\pi)^3 |\zeta(k_1)|^2 \delta^3(\mathbf{k}_1 - \mathbf{k}_2)$$

IS CONSTANT AFTER HORIZON EXIT. (LOOPS $\sim N_e(t)$)

[WEINBERG 2006]

- **WHAT ABOUT HIGHER POINT FUNCTIONS ?**

CONSTANCY CONDITION FOR $\zeta(k)$ 3-POINT FUNCTION

TO OBTAIN $\langle \hat{\zeta}^n \rangle$ USE THE IN-IN FORMALISM TO OBTAIN CORRELATIONS AS A FUNCTION OF TIME

$$\langle \hat{\zeta}(t)^n \rangle = \sum_{N=0}^{\infty} i^N \int_{t_0}^t dt_N \int_{t_0}^{t_N} dt_{N-1} \cdots \int_{t_0}^{t_2} dt_1 \times \left\langle \left[\hat{H}_I(t_1), \left[\hat{H}_I(t_2), \cdots \left[\hat{H}_I(t_N), \hat{\zeta}_I^n(t) \right] \cdots \right] \right] \right\rangle$$

\hat{H}_I IS THE INTERACTION HAMILTONIAN

$\hat{\zeta}_I$ = INTERACTION PICTURE, $\langle \dots \rangle = {}_{\text{in}} \langle 0 | \dots | 0 \rangle_{\text{in}}$

$$\langle \hat{\zeta}(t)^3 \rangle = i \int_{t_0}^t dt' \left\langle \left[\hat{H}_I(t'), \hat{\zeta}_I(t)^3 \right] \right\rangle \quad \text{LOWEST ORDER}$$

CONSTANCY CONDITION FOR $\zeta(k)$ 3-POINT FUNCTION

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- FOR $\langle \hat{\zeta}(t)^3 \rangle$ TO BE CONSTANT AFTER HORIZON EXIT THE CONTRIBUTION TO THE INTEGRAL FOR t AFTER HORIZON EXIT SHOULD BE SUPPRESSED.

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- WEINBERG (2008) SHOWED THIS WAS SO, BUT .. FOR GAUSSIAN INFLATON FLUCTUATIONS, I.E., IGNORING SELF INTERACTIONS OF THE INFLATON.

CONSTANCY CONDITION FOR $\zeta(k)$ 3-POINT FUNCTION

- ζ IS CLASSICAL AND CONSTANCY OUTSIDE THE HORIZON IMPLIES CORRELATIONS ALSO CONSTANT OUTSIDE THE HORIZON
- WEINBERG (2006) ARGUES THAT PERTURBATIONS OUTSIDE HORIZON CLASSICAL IN THE SENSE THAT

$$[\hat{\zeta}, \hat{\zeta}], [\hat{\zeta}, \dot{\hat{\zeta}}] \rightarrow 0 \quad \text{for large } t$$

n-POINT FUNCTIONS OF CAN GROW AS $(\ln a)$ OR POWERS OF $(\ln a)$, BUT NOT POWERS OF a .

$$(\ln a) = N_e$$

CONSTANCY CONDITION FOR $\zeta(k)$ 3-POINT FUNCTION

- WE INVESTIGATE WHETHER THE 3-POINT FUNCTION IS CONSTANT OUTSIDE THE HORIZON, IN LIGHT OF THE TIME DEPENDENT CONTRIBUTION PROPORTIONAL TO N_e FROM INFLATON SELF INTERACTIONS

CALCULATION OF THE 3-POINT FUNCTION OF $\zeta(k)$ AND f_{NL}

- WE CALCULATE $\langle (\hat{\delta\phi})^3 \rangle$ IN THE $\delta\phi \neq 0$ GAUGE, USING THE CANONICAL FORMALISM FOR CUBIC NEW INFLATION $V(\phi) = V_0 - \mu \phi^3$.
- WE THEN RELATE $\zeta(k,t)$ TO $\delta\phi(k,t)$. AND CALCULATE $\langle \hat{\zeta}^3(t) \rangle$, AND THE NON-GAUSSIANITY PARAMETER f_{NL} .
- t IS ARBITRARY, UNLIKE IN THE δN FORMALISM WHERE $t \sim$ TIME OF HORIZON EXIT SO AS TO STUDY POSSIBLE GROWTH AFTER HORIZON EXIT

CALCULATION OF THE 3-POINT FUNCTION OF $\zeta(k)$ AND f_{NL}

$\langle (\hat{\delta\phi})^3 \rangle$ IN THE $\delta\phi \neq 0$ GAUGE, FOR CUBIC NEW
INFLATION $V(\phi) = V_0 - \mu \phi^3$.

$$\langle \hat{\varphi}(t)^3 \rangle = i \int_{t_0}^t dt' \left\langle \left[\hat{H}_I(t'), \hat{\varphi}_I(t)^3 \right] \right\rangle$$

$$\varphi \equiv \delta\phi \text{ AND } \phi(x, t) = \phi(t) + \varphi(x, t)$$

TO OBTAIN H_I WITH CUBIC INTERACTION TERMS

$$S[\varphi, \dot{\varphi}] = S_2 + S_3$$

CALCULATION OF THE 3-POINT FUNCTION OF $\zeta(k)$ AND f_{NL}

$\langle (\hat{\delta}\phi)^3 \rangle$ IN THE $\delta\phi \neq 0$ GAUGE, FOR CUBIC NEW
INFLATION $V(\phi) = V_0 - \mu \phi^3$.

$$\langle \hat{\varphi}(t)^3 \rangle = i \int_{t_0}^t dt' \left\langle \left[\hat{H}_I(t'), \hat{\varphi}_I(t)^3 \right] \right\rangle \quad \varphi \equiv \delta\phi$$

$$S_3 = \int dt d^3x a^3 \left[-\frac{\dot{\varphi}}{4H} \varphi \dot{\varphi}^2 - \frac{1}{a^2} \frac{\dot{\phi}}{4H} \varphi (\partial_i \varphi)^2 \right. \\ \left. + \frac{\dot{\phi}}{2H} \partial_i \varphi (\partial_i^{-1} \dot{\varphi}) \dot{\varphi} - \frac{1}{6} V'''(\phi) \varphi^3 \right]$$

CALCULATION OF THE 3-POINT FUNCTION OF $\zeta(k)$ AND f_{NL}

$$\langle \hat{\varphi}(t)^3 \rangle = i \int_{t_0}^t dt' \left\langle \left[\hat{H}_I(t'), \hat{\varphi}_I(t)^3 \right] \right\rangle$$

1. $H_I = \frac{\dot{\varphi}}{4H} \varphi \dot{\varphi}^2$

$$I_1 = -2 \times \frac{1}{4} \frac{H^3(\tau)}{\prod_i (2k_i^3)} (2\pi)^3 \delta^3 \left(\sum_i \vec{k}_i \right) \dot{\phi}(\tau) \left[\frac{k_2^2 k_3^2}{k_t} + \frac{k_1 k_2^2 k_3^2}{k_t^2} + \text{perm} \right]$$

τ IS THE CONFORMAL TIME

CALCULATION OF THE 3-POINT FUNCTION OF $\zeta(k)$ AND f_{NL}

$$\langle \hat{\varphi}(t)^3 \rangle = i \int_{t_0}^t dt' \left\langle \left[\hat{H}_I(t'), \hat{\varphi}_I(t)^3 \right] \right\rangle$$

$$2. H_I = \frac{1}{a^2} \frac{\dot{\phi}}{4H} \varphi (\partial_i \varphi)^2$$

$$I_2 = -2 \frac{1}{4} \frac{H^3(\tau)}{\prod_i (2k_i^3)} (2\pi)^3 \delta^3 \left(\sum_i \vec{k}_i \right) \dot{\phi}(\tau) (\vec{k}_2 \cdot \vec{k}_3) \\ \times \left[-k_t + \sum_{i \neq j} \frac{k_i k_j}{k_t} + \frac{k_1 k_2 k_3}{k_t^2} + \text{perm} \right]$$

CALCULATION OF THE 3-POINT FUNCTION OF $\zeta(k)$ AND f_{NL}

$$\langle \hat{\varphi}(t)^3 \rangle = i \int_{t_0}^t dt' \left\langle \left[\hat{H}_I(t'), \hat{\varphi}_I(t)^3 \right] \right\rangle$$

$$3. H_I = -\frac{\dot{\phi}}{2H} \partial_i \varphi (\partial_i^{-1} \dot{\varphi}) \dot{\varphi}$$

$$I_3 \equiv 2 \frac{1}{2} \frac{H^3(\tau)}{\prod_i (2k_i^3)} (2\pi)^3 \delta^3 \left(\sum_i \vec{k}_i \right) \dot{\phi}(\tau) (\vec{k}_1 \cdot \vec{k}_2) \left[\frac{k_3^2}{k_t} + \frac{k_1 k_3^2}{k_t^2} + \text{perm} \right]$$

CALCULATION OF THE 3-POINT FUNCTION OF $\zeta(k)$ AND f_{NL}

$$\langle \hat{\varphi}(t)^3 \rangle = i \int_{t_0}^t dt' \left\langle \left[\hat{H}_I(t'), \hat{\varphi}_I(t)^3 \right] \right\rangle$$

$$4. H_I = \frac{1}{6} V'''(\phi) \varphi^3$$

$$I_4 = \frac{H^2 V'''}{4 \prod_i k_i^3} (2\pi)^3 \delta^3 \left(\sum_i \vec{k}_i \right) \\ \times \left[-\frac{4}{9} k_t^3 + k_t \sum_{i < j} k_i k_j + \frac{1}{3} \left\{ \frac{1}{3} + \gamma + \ln |k_t \tau| \right\} \sum_i k_i^3 \right]$$

IS EULER'S CONSTANT = 0.577216

CALCULATION OF THE 3-POINT FUNCTION OF $\zeta(\mathbf{k})$ AND f_{NL}

$$\langle \hat{\varphi}(\vec{k}_1, t) \hat{\varphi}(\vec{k}_2, t) \hat{\varphi}(\vec{k}_3, t) \rangle = (2\pi)^3 \delta(\vec{k}_1 + \vec{k}_2 + \vec{k}_3) \\ \times \left[\frac{H^2 V''''}{4 \prod_i k_i^3} \left(-\frac{4}{9} k_t^3 + k_t \sum_{i < j} k_i k_j + \frac{1}{3} \left\{ \frac{1}{3} + \gamma + \ln |k_t \tau| \right\} \sum_i k_i^3 \right) \right. \\ \left. + \frac{H^4}{8 \prod_i k_i^3} \frac{\dot{\phi}}{H} \left(\frac{1}{2} \sum_i k_i^3 - \frac{4}{k_t} \sum_{i < j} k_i^2 k_j^2 - \frac{1}{2} \sum_{i \neq j} k_i k_j^2 \right) \right]$$

$$i, j = 1, 2, 3, \quad k_i = |\mathbf{k}_i| \quad k_t = \sum k_i, \quad k_i \text{ APPROX EQUAL}$$

SEERY, LIDSEY (PI w/o SELF INTER), SEERY ET AL (HEIS. FIELD EQS)

CALCULATION OF THE 3-POINT FUNCTION OF $\zeta(\mathbf{k})$ AND f_{NL}

$$\zeta(\mathbf{k}, t) = -\frac{1}{\sqrt{2\epsilon}}\varphi(\mathbf{k}, t) + \frac{1}{2} \left(1 - \frac{\eta}{2\epsilon}\right) \int \frac{d^3q}{(2\pi)^3} \varphi(\mathbf{k}_1 - \mathbf{q}, t)\varphi(\mathbf{q}, t) + \dots$$

[GAUGE TRANSFORMATION - MALDACENA 2003]

$$\langle \hat{\zeta}(\mathbf{k}_1, t) \hat{\zeta}(\mathbf{k}_2, t) \hat{\zeta}(\mathbf{k}_3, t) \rangle = -\frac{1}{(2\epsilon)^{\frac{3}{2}}} \langle \hat{\varphi}(\mathbf{k}_1, t) \hat{\varphi}(\mathbf{k}_2, t) \hat{\varphi}(\mathbf{k}_3, t) \rangle + \frac{1}{2\epsilon} \frac{1}{2} \left(1 - \frac{\eta}{2\epsilon}\right) \langle \hat{\varphi}(\mathbf{k}_1, t) \hat{\varphi}(\mathbf{k}_2, t) \int \frac{d^3q}{2\pi^3} \hat{\varphi}(\mathbf{k}_3 - \mathbf{q}, t) \hat{\varphi}(\mathbf{q}, t) + \text{perm} \rangle$$

CALCULATION OF THE 3-POINT FUNCTION OF $\zeta(\mathbf{k})$ AND f_{NL}

$$\langle \hat{\zeta}(\mathbf{k}_1, t) \hat{\zeta}(\mathbf{k}_2, t) \hat{\zeta}(\mathbf{k}_3, t) \rangle \equiv (2\pi)^3 \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3) \frac{6}{5} f_{NL} \sum_{i < j} P_{\zeta}(k_i) P_{\zeta}(k_j)$$

$$i, j = 1, 2, 3, \quad k_i = |\mathbf{k}_i|$$

CURVATURE POWER SPECTRUM $P_{\zeta}(k) = \frac{1}{2\epsilon} \frac{H^2}{2k^3}$

$$\langle \hat{\zeta}(\mathbf{k}) \hat{\zeta}(\mathbf{k}') \rangle = (2\pi)^3 \delta^3(\mathbf{k} + \mathbf{k}') P_{\zeta}(k)$$

$$\left[(k^3 / 2\pi^2) P_{\zeta}(k) \right]^{\frac{1}{2}} = 5 \times 10^{-5}$$

CALCULATION OF THE 3-POINT FUNCTION OF $\zeta(k)$ AND f_{NL}

$$\frac{6}{5} f_{\text{NL}}(k_1, k_2, k_3, t) = \underline{\xi} \left[\frac{1}{3} + \gamma - \underline{N_e} + \frac{3}{\sum_i k_i^3} \left(k_t \sum_{i<j} k_i k_j - \frac{4}{9} k_t^3 \right) \right]$$

$$+ \frac{3}{2} \epsilon - \underline{\eta} + \frac{\epsilon}{\sum_i k_i^3} \left(\frac{4}{k_t} \sum_{i<j} k_i^2 k_j^2 + \frac{1}{2} \sum_{i \neq j} k_i k_j^2 \right)$$

$$k_t = \sum k_i, \quad k_i \text{ APPROX EQUAL}$$

FOR $n_s = 0.96$, $\eta = -0.02$. $\xi = 0.5 \eta^2$. $\xi N_e = 0.012$, $\epsilon \ll \xi N_e$

SO THE SELF INTERACTION CONTRIBUTION $\sim \xi$ SHOULD
NOT BE IGNORED OUTRIGHT

$d f_{\text{NL}} / dt$

- f_{NL} IS A FUNCTION OF $N_e(t) = H(t - t_{\text{exit}})$, AND $\epsilon(t)$, $\eta(t)$, $\xi(t)$.
- DOMINANT CONTRIBUTION TO f_{NL} IS

$$f_{\text{NL}} \approx \frac{5}{6} [-\xi N_e - \eta]$$

$$\frac{d}{dt} f_{\text{NL}} \approx \frac{5}{6} \frac{d}{dt} [-\xi N_e - \eta]$$

$d f_{\text{NL}} / dt$

- f_{NL} IS A FUNCTION OF $N_e(t) = H (t - t_{\text{exit}})$, AND $\epsilon(t)$, $\eta(t)$, $\xi(t)$.

$$\frac{d}{dt} f_{\text{NL}} \approx \frac{5}{6} \frac{d}{dt} [-\xi N_e - \eta]$$

$$\frac{d\eta}{dt} \simeq [2\epsilon\eta - \xi] H \quad \text{AND} \quad \frac{d\xi}{dt} \simeq [4\epsilon\xi - \eta\xi] H$$

$$\frac{d}{dt} f_{\text{NL}} = \frac{5}{6} [-\xi H + \xi H]$$

$$= 0$$

SIMILAR TO SEERY, MALIK, LYTH (2008)

$$d f_{NL} / dt = 0$$

TIME EVOLUTION OF f_{NL} DUE TO SELF INTERACTIONS IS CANCELLED BY CONTRIBUTION OF OTHER TERMS FROM COSMOLOGICAL PERT. THEORY.

THUS THE TOTAL 3-POINT FUNCTION OF ζ DOES NOT GROW OUTSIDE THE HORIZON.

CONCLUSION

- THE ARGUMENT THAT ONE SHOULD IGNORE THE CONTRIBUTION OF SELF INTERACTIONS TO f_{NL} HAS BEEN RE-EXAMINED
- GROWS OUTSIDE THE HORIZON $\sim N_e$
- WE CALCULATE $\langle \hat{\zeta}^3(t) \rangle$ TO STUDY SUPERHORIZON EVOLUTION
- GROWTH DUE TO SELF INTERACTIONS IS CANCELLED BY GROWTH IN OTHER TERMS IN f_{NL}

CONCLUSION

- THE ARGUMENT THAT ONE SHOULD IGNORE THE CONTRIBUTION OF SELF INTERACTIONS TO f_{NL} HAS BEEN RE-EXAMINED
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THANK YOU